

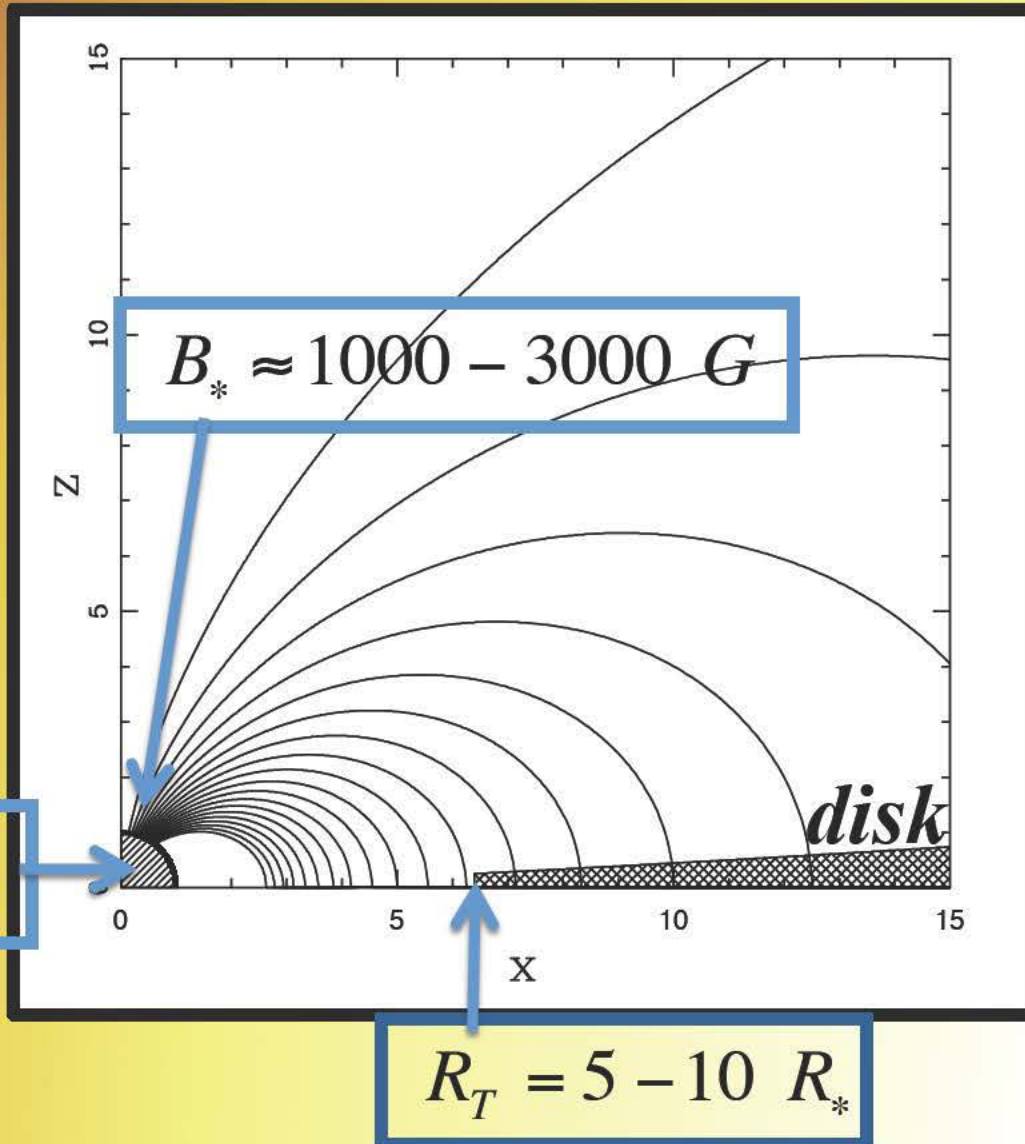
Magnetically Controlled Accretion Flows onto Young Stellar Objects

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(w: S. Gregory, CalTech)***

Motivation

- Most material that ends up in a forming star is processed through the disk; final accretion stage occurs via magnetic field
- T Tauri stars are observed to have magnetic fields with strong Octupole Components (in addition to Dipoles)
- Want to understand the transition through sonic points for magnetically controlled flows in general

The Star/Disk System



Basic Regime of Operation

$$\frac{dM}{dt} \approx 10^{-7} - 10^{-8} M_{sun} yr^{-1}$$

$$\frac{B^2}{8\pi\rho v^2} \approx 350 - 3000 \text{ (magnetically - controlled)}$$

$$\frac{\omega_c}{\Gamma} = \frac{qB}{cmn\sigma} \approx 10^4 - 10^5 \text{ (well - coupled)}$$

$$\frac{B_{\perp}}{B} = O(8\pi\rho v^2 / B^2) < 10^{-3} \text{ (current - free)}$$

Equations of Motion

Steady-state flow, polytropic equation of state:

$$\nabla \cdot (\rho \vec{u}) = 0 \qquad P = K \rho^{1+1/n}$$

$$\vec{u} \cdot \nabla \vec{u} + \nabla \Psi + \frac{1}{\rho} \nabla P + \vec{\Omega} \times (\vec{\Omega} \times \vec{r}) = 0$$

$$\vec{B} = \kappa \rho \vec{u} \quad \text{where} \quad \kappa = \text{const}$$

Construct Coordinate Systems that follow Magnetic Field Lines

Basis Vectors

covariant :

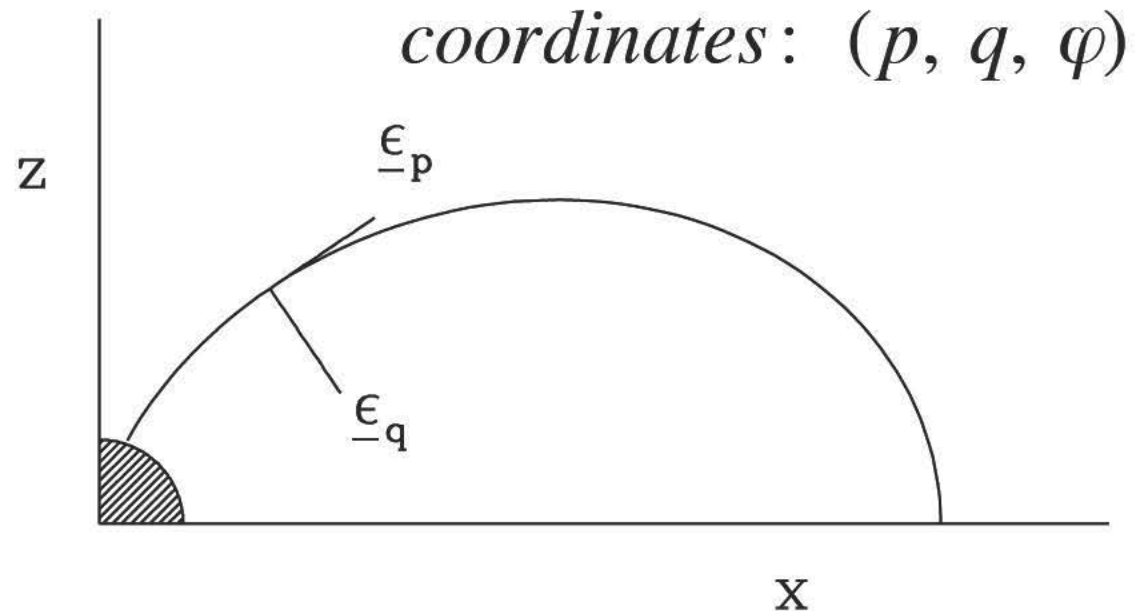
$$\underline{\varepsilon}_p = \nabla p$$

$$\underline{\varepsilon}_q = \nabla q$$

contravariant :

$$\vec{e}_p = \partial \vec{r} / \partial p$$

$$\vec{e}_q = \partial \vec{r} / \partial q$$



Scale Factors: $h_j = |\underline{\varepsilon}_j|^{-1}$

Unit Vectors: $\hat{e}_j = \underline{\varepsilon}_j h_j = \vec{e}_j h_j^{-1}$

Dimensionless Equations of Motion

Steady-state flow along field-line direction:
Fluid fields are functions of coordinate p only.

$$\alpha \frac{\partial u}{\partial p} + u \frac{\partial \alpha}{\partial p} = - \frac{\alpha u}{h_q h_\phi} \frac{\partial}{\partial p} (h_q h_\phi)$$

$$u \frac{\partial u}{\partial p} + \frac{\alpha^{1/n}}{\alpha} \frac{\partial \alpha}{\partial p} - \omega \xi \sin \theta |\nabla p|^{-1} (\hat{x} \cdot \hat{p}) = - \frac{\partial \psi}{\partial p}$$

$$u = \frac{|\vec{u}|}{a_s}, \quad \alpha = \frac{\rho}{\rho_1}, \quad \psi = \frac{\Psi}{a_s^2}, \quad b \equiv \frac{GM_P}{R_* a_s^2} \quad \omega \equiv \left(\frac{\Omega R_*}{a_s} \right)^2$$

Integrated Equations of Motion

$$\alpha u h_q h_\phi = \lambda$$

$$\frac{1}{2} u^2 + n \alpha^{1/n} + \psi - \omega I = \varepsilon$$

$$I = \int \xi \sin \theta |\nabla p|^{-1} (\hat{x} \cdot \hat{p}) dp$$

$\lambda = \text{mass accretion rate}$

$\varepsilon = \text{energy}$

Sonic Transition Condition

$$\alpha^{1/n} \frac{Y}{\xi} + \omega \Lambda = \frac{b}{\xi^2}$$

$$Y(\xi, \theta) \equiv \frac{\xi}{h_p} \frac{\partial h_p}{\partial \xi} \quad (\text{where } h_p = h_q h_\phi)$$

$$\Lambda(\xi, \theta) \equiv \xi \sin \theta \left(\frac{\partial p}{\partial \xi} \right)^{-1} (\hat{x} \cdot \nabla p)$$

General Constraint on Steady Polytropic Transonic Accretion Flow

Observations show that flow must be *transonic*. *Steady-state* accretion solutions that pass through the sonic point and approach free-fall speed near the star must satisfy the constraint:

$$n > \ell + 3/2 \rightarrow 9/2 \text{ (octupole)}$$

Steady flow must be nearly isothermal for fields with higher order multipoles.

Score Card

b and ω : system parameters ($n \rightarrow \infty$)

λ and ε : conserved quantities

$Y(\xi, \theta)$ and $\Lambda(\xi, \theta)$: functions that
specify magnetic field geometry

field lines $\Rightarrow q = \text{const} \Rightarrow \theta = F(\xi)$

Dipole Coordinate System

$$\vec{B} = B_0 \left[\xi^{-3} (3 \cos \theta \hat{r} - \hat{z}) \right] \quad \text{where} \quad \xi = r / R_*$$

$$p = -\xi^{-2} \cos \theta \quad \text{and} \quad q = \xi^{-1} \sin^2 \theta$$

$$\nabla p = 2\xi^{-3} \cos \theta \hat{r} + \xi^{-3} \sin \theta \hat{\theta}$$

$$\nabla q = -\xi^{-2} \sin^2 \theta \hat{r} + 2\xi^{-2} \cos \theta \sin \theta \hat{\theta}$$

$$\left. \begin{array}{l} \nabla p = 2\xi^{-3} \cos \theta \hat{r} + \xi^{-3} \sin \theta \hat{\theta} \\ \nabla q = -\xi^{-2} \sin^2 \theta \hat{r} + 2\xi^{-2} \cos \theta \sin \theta \hat{\theta} \end{array} \right\} \nabla p \cdot \nabla q = 0$$

$$h_p = \xi^3 \left[4 \cos^2 \theta + \sin^2 \theta \right]^{-1/2}$$

$$h_p = \frac{\xi^2}{\sin \theta} \left[4 \cos^2 \theta + \sin^2 \theta \right]^{-1/2} \quad \left(h_\phi = \xi \sin \theta \right)$$

Ancillary Functions

For Coordinate System that follows
Dipole Magnetic Field Lines:

$$Y = Y(\xi) = \frac{3}{2} \frac{8 - 5q\xi}{4 - 3q\xi}$$

$$\Lambda = \Lambda(\xi) = \frac{3}{2} q \xi^2$$

*along field line
labeled by the
coordinate q*

Solutions (Isothermal Limit)

Accretion flow follows magnetic field lines, which are lines of constant coordinate q .

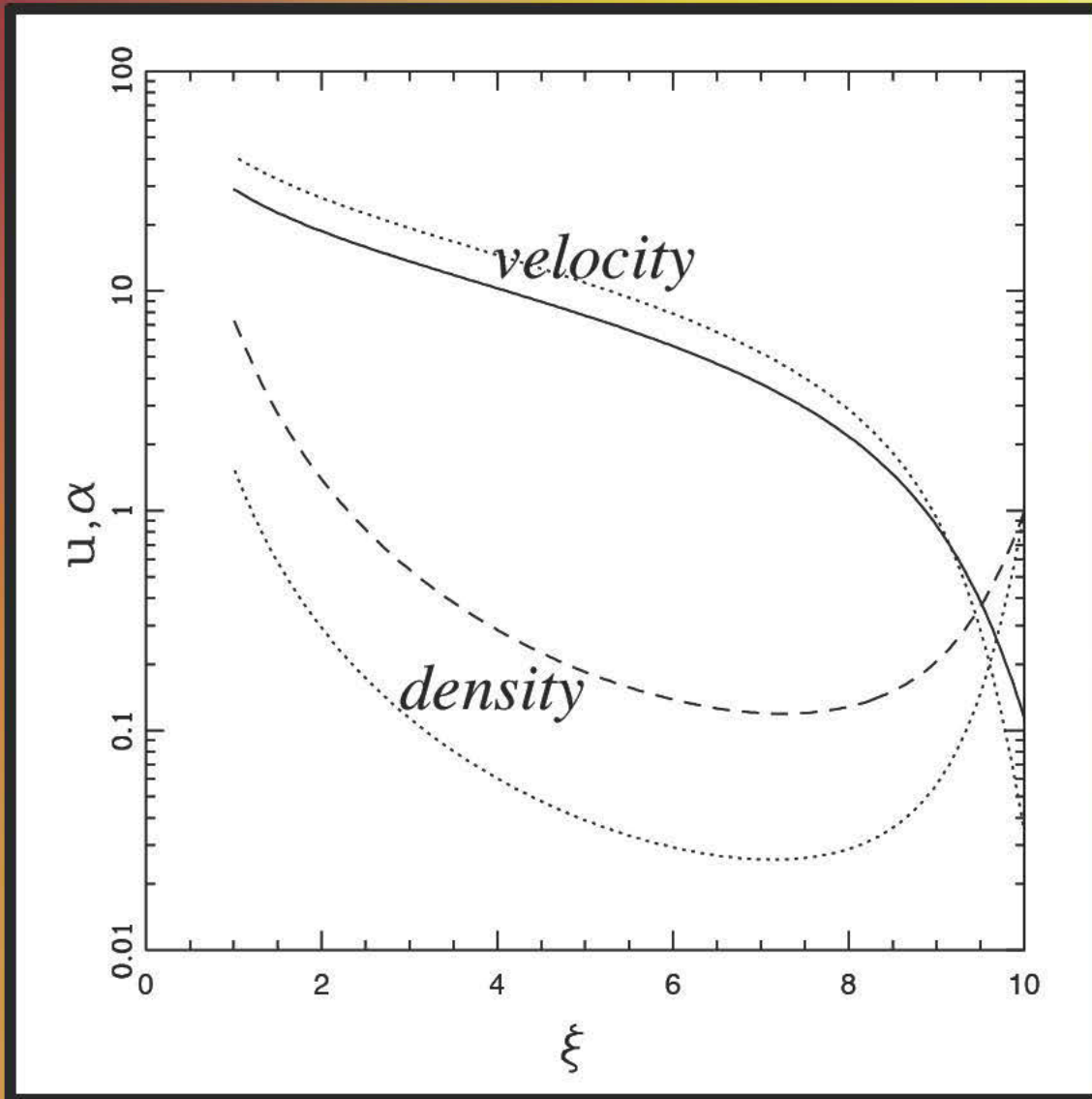
$$3\xi \frac{8 - 5q\xi}{4 - 3q\xi} = 2b - 3 \omega q \xi^4 \quad \text{(Sonic Point Condition)}$$

$$\log \lambda - \frac{1}{2} \lambda^2 = 3 \log \xi_s - \frac{1}{2} \log(4 - 3\xi_s) - \frac{1}{2}$$

(Mass Accretion Rate)

$$+ b \left(\frac{1}{\xi_s} + \frac{1}{2} \xi_s^2 - \frac{3}{2} \right)$$

Dipole Accretion Solution



dipole fields

isothermal flow

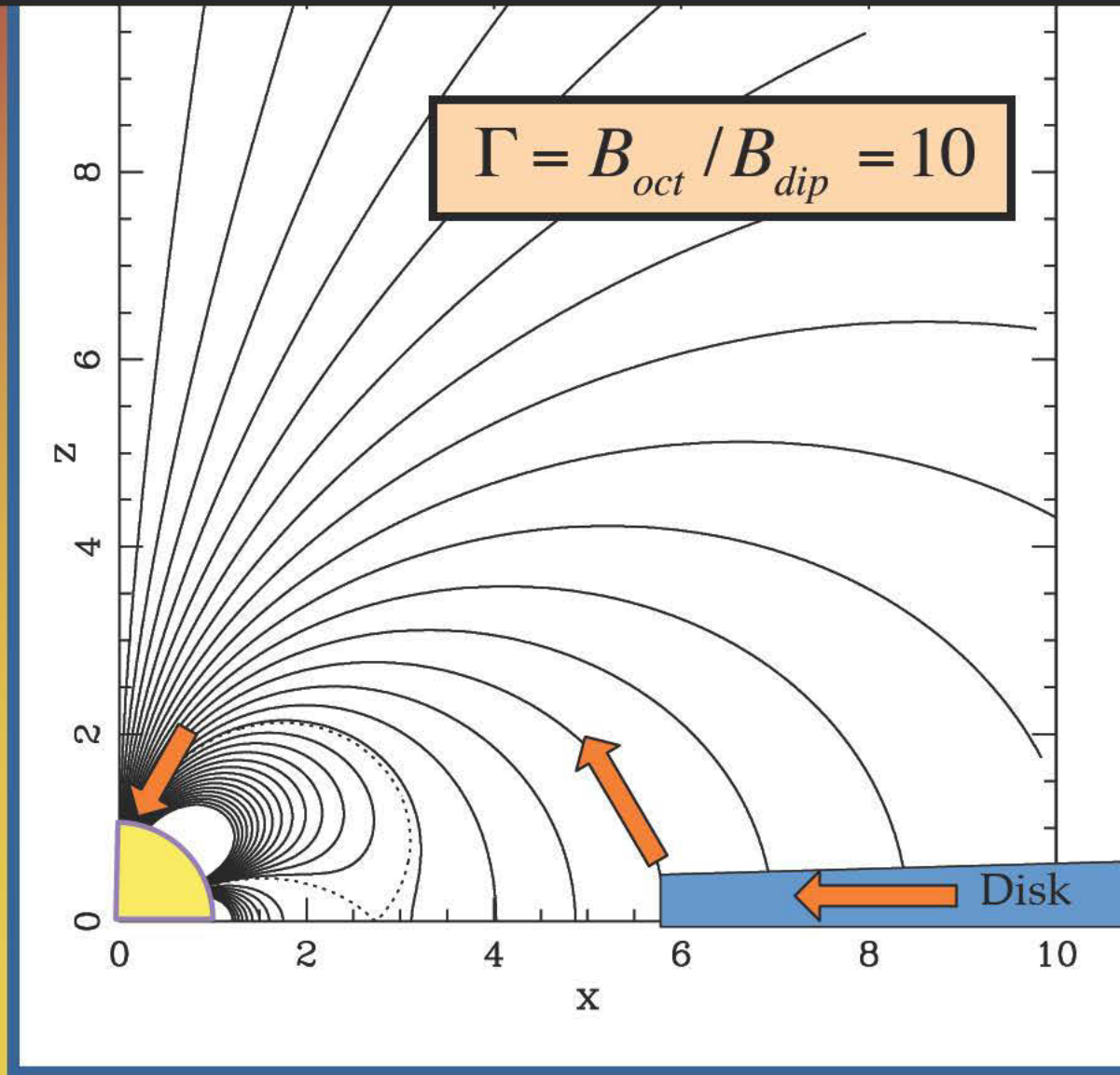
$$\xi_d = 10$$

$$\xi_* = 1$$

$$b = 500$$

$$b = 1000 \text{ (dots)}$$

Dipole + Octupole Configuration



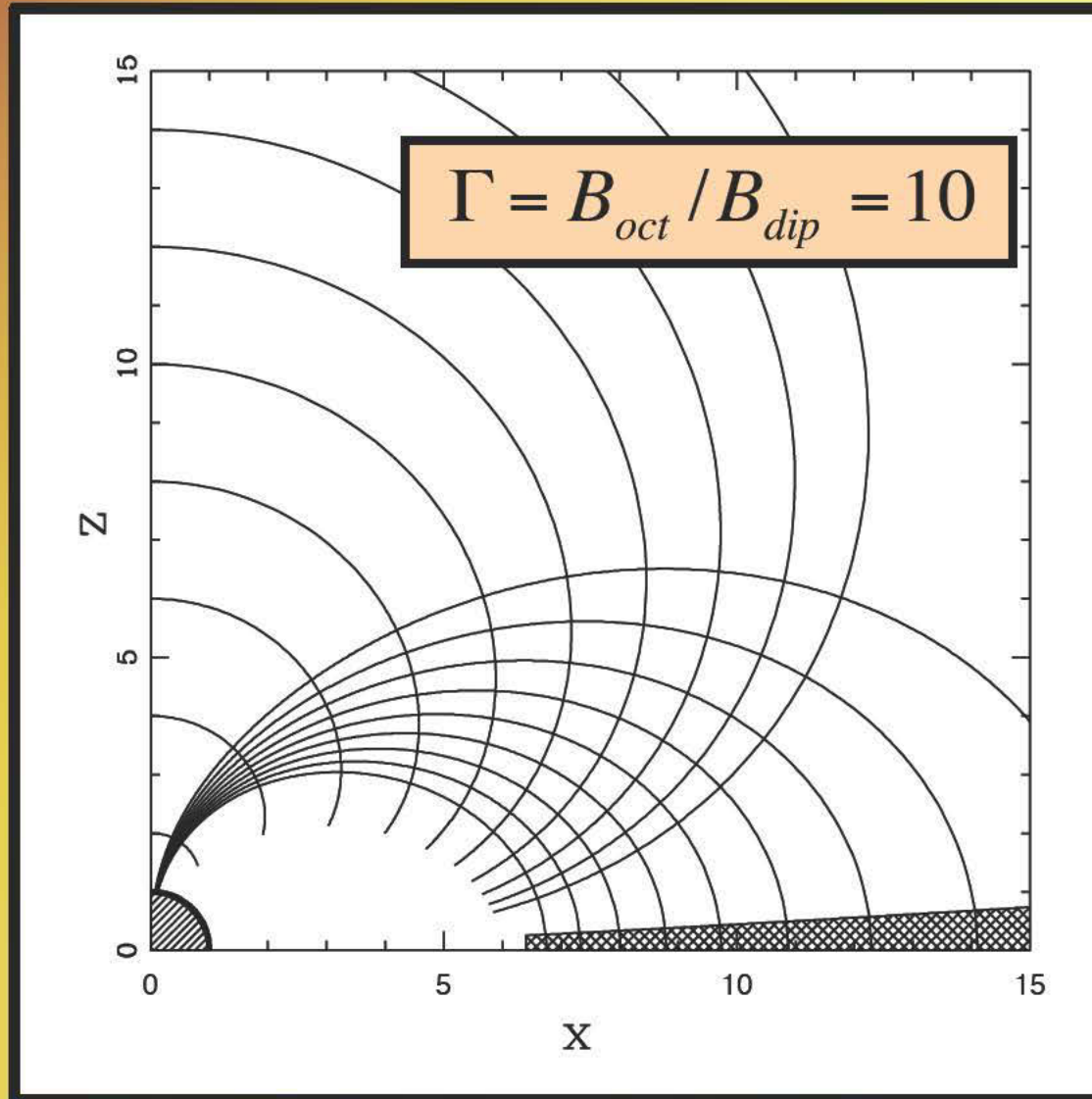
Dipole + Octupole Coordinate System

$$\vec{B} = B_{dip} \xi^{-3} \left(2 \cos \theta \hat{r} + \sin \theta \hat{\theta} \right) +$$

$$\frac{1}{2} B_{oct} \xi^{-5} \left[(5 \cos^2 \theta - 3) \cos \theta \hat{r} + \frac{3}{4} (5 \cos^2 \theta - 1) \sin \theta \hat{\theta} \right]$$

$$\left. \begin{aligned} p &= -\frac{\Gamma}{4} \xi^{-4} (5 \cos^2 \theta - 3) \cos \theta - \xi^{-2} \cos \theta \\ q &= \frac{\Gamma}{4} \xi^{-3} (5 \cos^2 \theta - 1) \sin^2 \theta + \xi^{-1} \sin^2 \theta \end{aligned} \right\} \text{where } \Gamma \equiv \frac{B_{oct}}{B_{dip}}$$

Dipole + Octupole Coordinate System



Dipole + Octupole Scale Factors

$$h_p = \xi^5 [f^2 \cos^2 \theta + g^2 \sin^2 \theta]^{-1/2}$$

$$h_q = \xi^4 (\sin \theta)^{-1} [f^2 \cos^2 \theta + g^2 \sin^2 \theta]^{-1/2}$$

where $f = \Gamma (5 \cos^2 \theta - 3) + 2\xi^2$

and $g = (3/4)\Gamma (5 \cos^2 \theta - 1) + \xi^2$

$$\sin^2 \theta = \frac{2}{5\Gamma} \left\{ (\xi^2 + \Gamma) - \left[(\xi^2 + \Gamma)^2 - 5\Gamma q \xi^3 \right]^{1/2} \right\}$$

so that $f, g, h_p, h_q = \text{Function}(\xi \text{ only})$

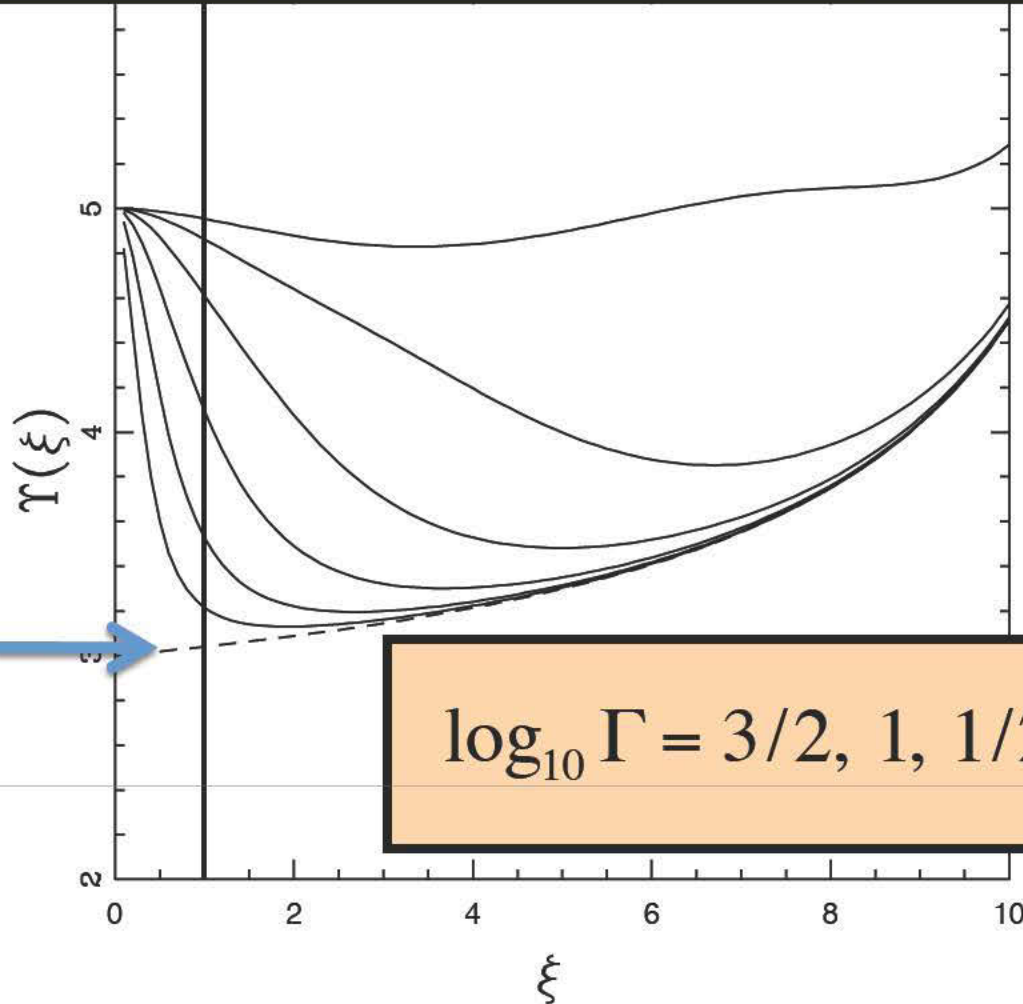
Ancillary Functions for Dipole + Octupole Configuration

$$\Lambda(\xi) = \frac{3\xi}{5\Gamma} \left\{ 1 + \frac{g(\xi)}{f(\xi)} \right\} \left\{ (\xi^2 + \Gamma) - [(\xi^2 + \Gamma)^2 - 5\Gamma q \xi^3]^{1/2} \right\}$$

$$Y(\xi) = 5 - \left[f^2 + (g^2 - f^2) \frac{1}{5\Gamma} (2\xi^2 + 2\Gamma - f) \right]^{-1} \frac{\xi}{5\Gamma} \times$$

$$\left\{ 5\Gamma f f_\xi + \left[g \left(\frac{3f_\xi}{4} - \xi \right) - f f_\xi \right] (2\xi^2 + 2\Gamma - f) + (g^2 - f^2) \left(2\xi - \frac{f_\xi}{2} \right) \right\}$$

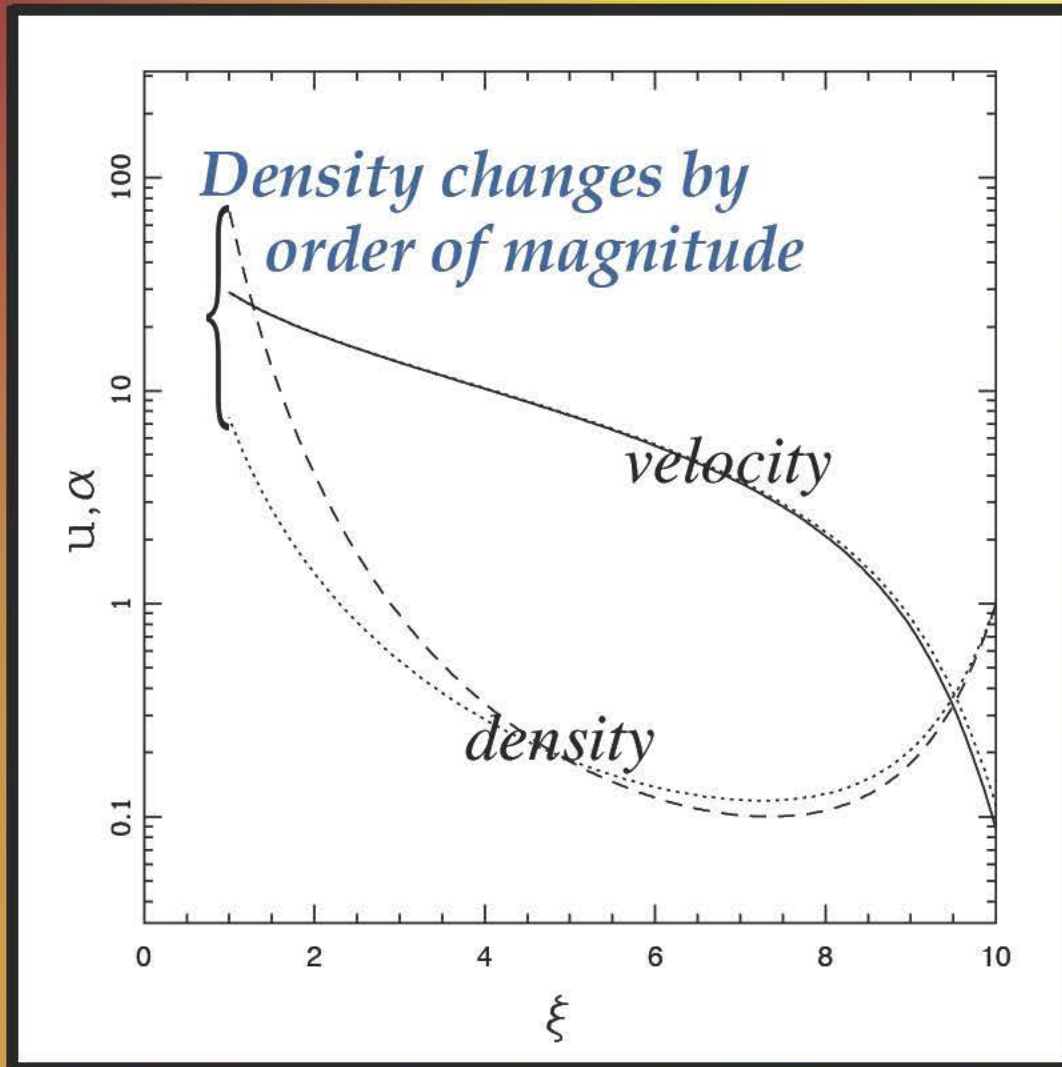
Index of Divergence Operator (Dipole + Octupole)



$\Gamma \rightarrow 0$

$$\log_{10} \Gamma = 3/2, 1, 1/2, 0, -1/2, -1$$

Accretion Solution (Dip+Oct)



isothermal flow

octupole $\Gamma = 10$

$$\xi_d = 10$$

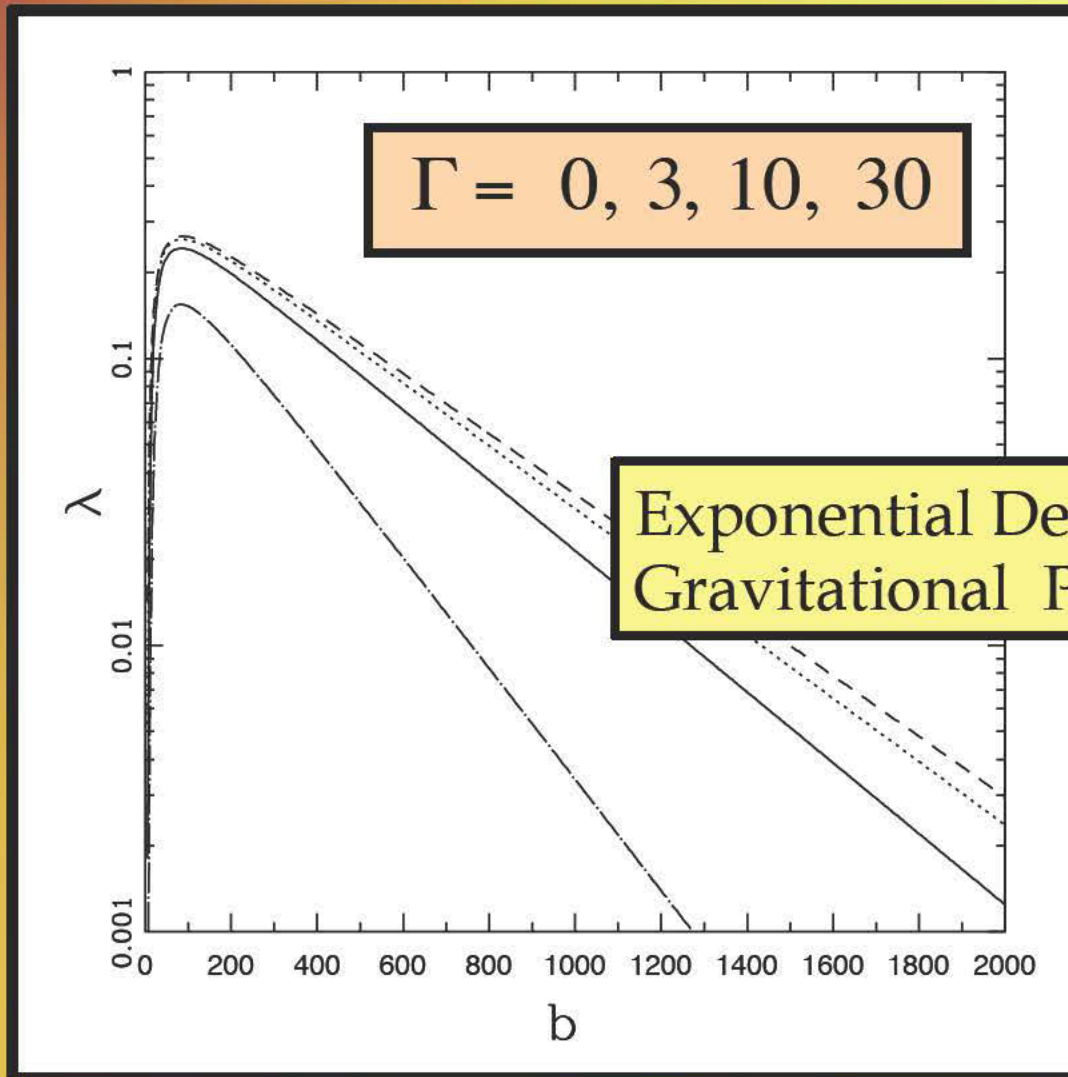
$$\xi_* = 1$$

$$b = 500$$

dots = dipole solution

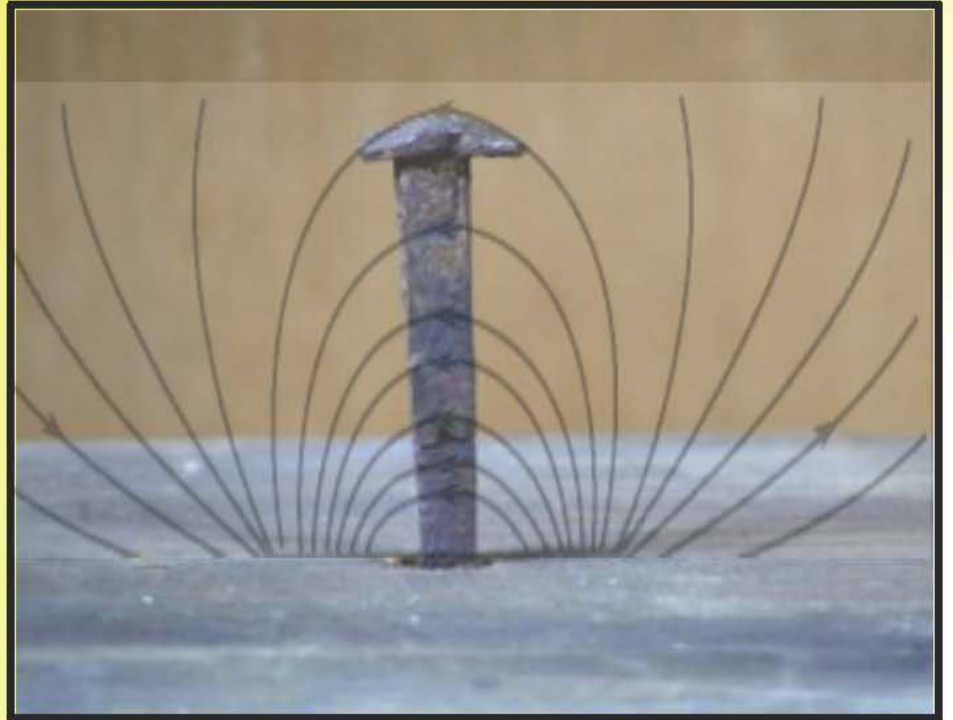
dashes = full solution

Dimensionless Accretion Rate

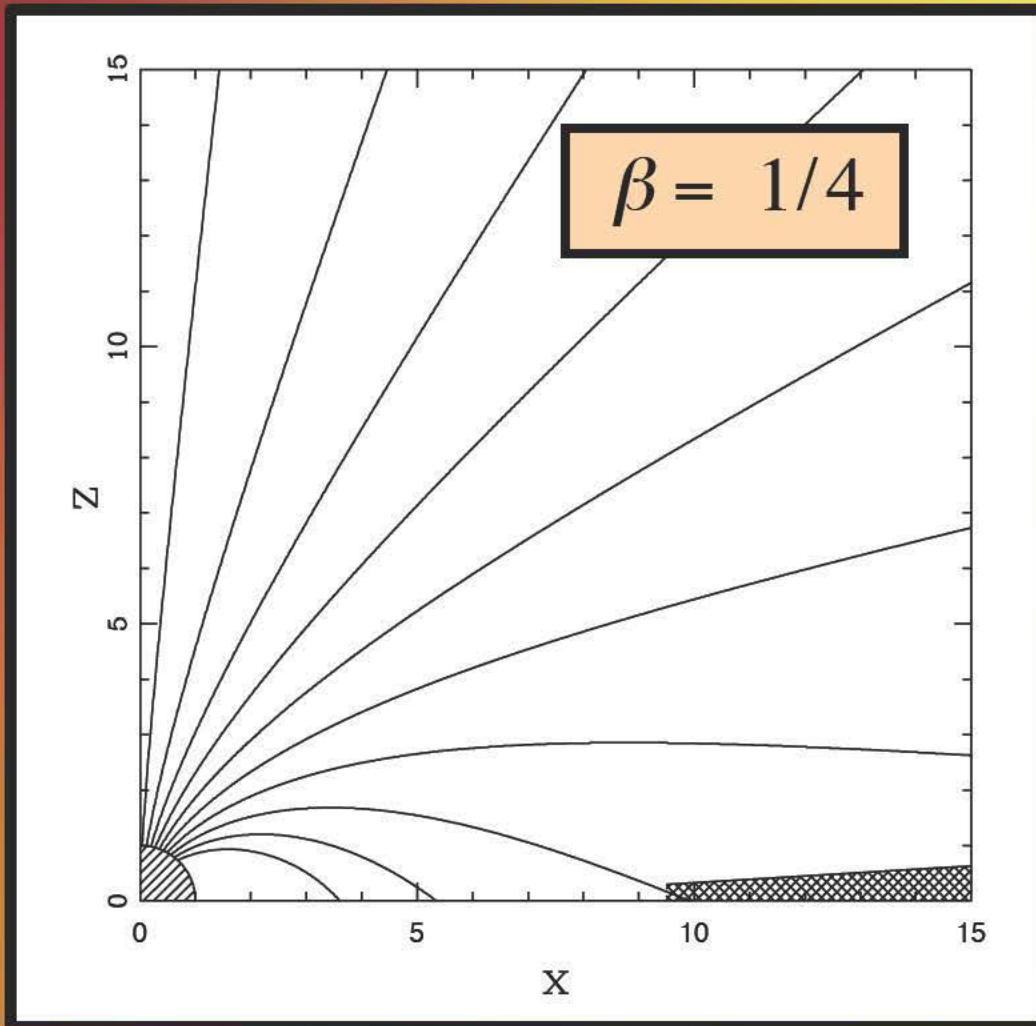




To a man with a hammer, everything looks like a nail



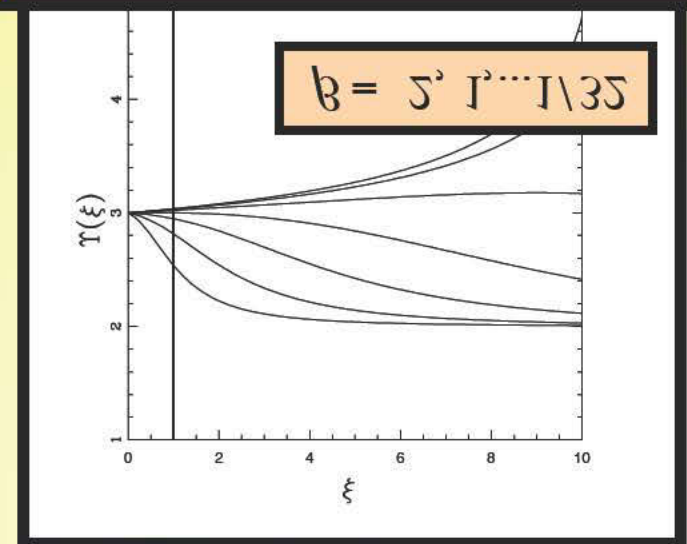
Dipole + Split-Monopole



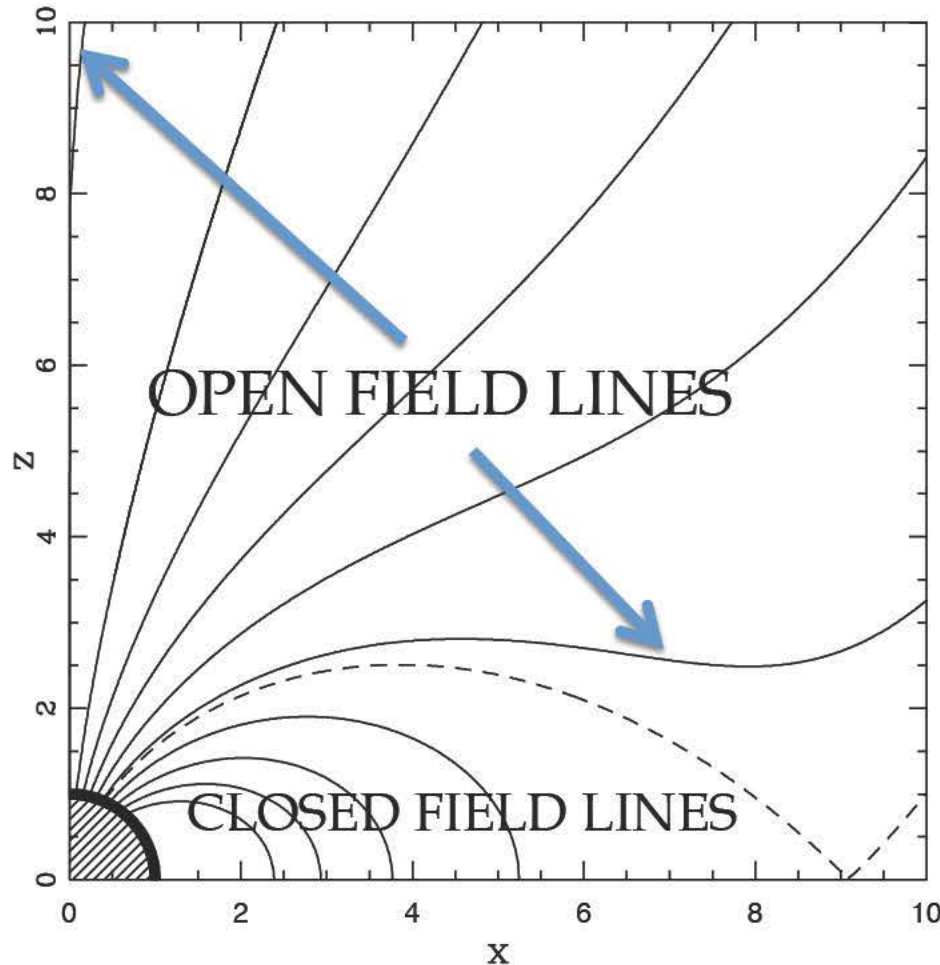
$$p = -\xi^{-2} \cos \theta - \beta \xi^{-1}$$

$$q = \xi^{-1} \sin^2 \theta - \beta \cos \theta$$

where $\beta \equiv 2B_{rad} / B_{dip}$



Planet Dipole in Stellar z-Field



$$p = (\beta\xi - \xi^{-2})\cos\theta$$

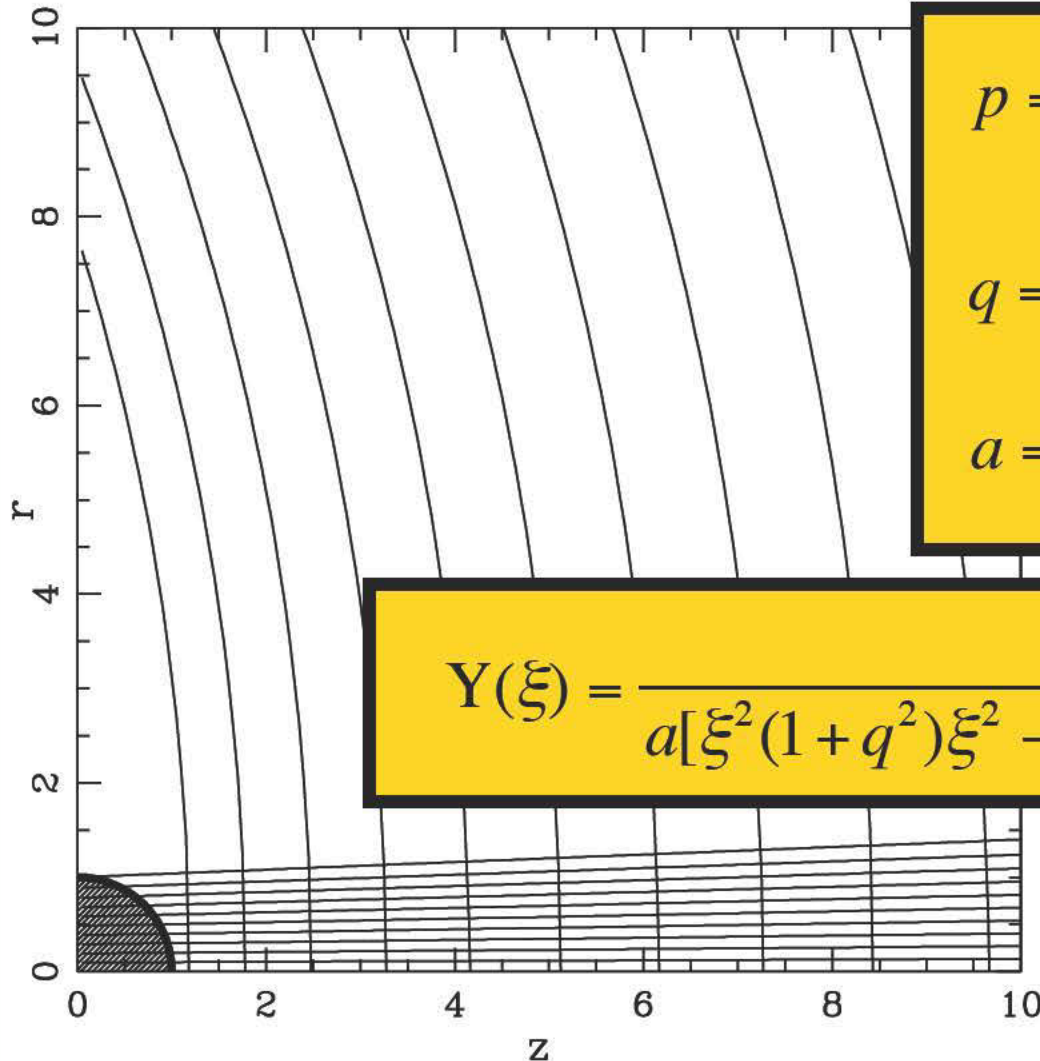
$$q = (\beta\xi^2 + 2/\xi)^{1/2}\sin\theta$$

$$\beta = (B_* R_*^3 / \varpi^3) / B_P$$

Field lines are open near planetary pole and are closed near the equator. Fraction of open field lines:

$$F_{open} = 1 - \left[1 - \frac{3\beta^{1/3}}{2 + \beta} \right]^{1/2}$$

Planet in Stellar Split-Monopole Field



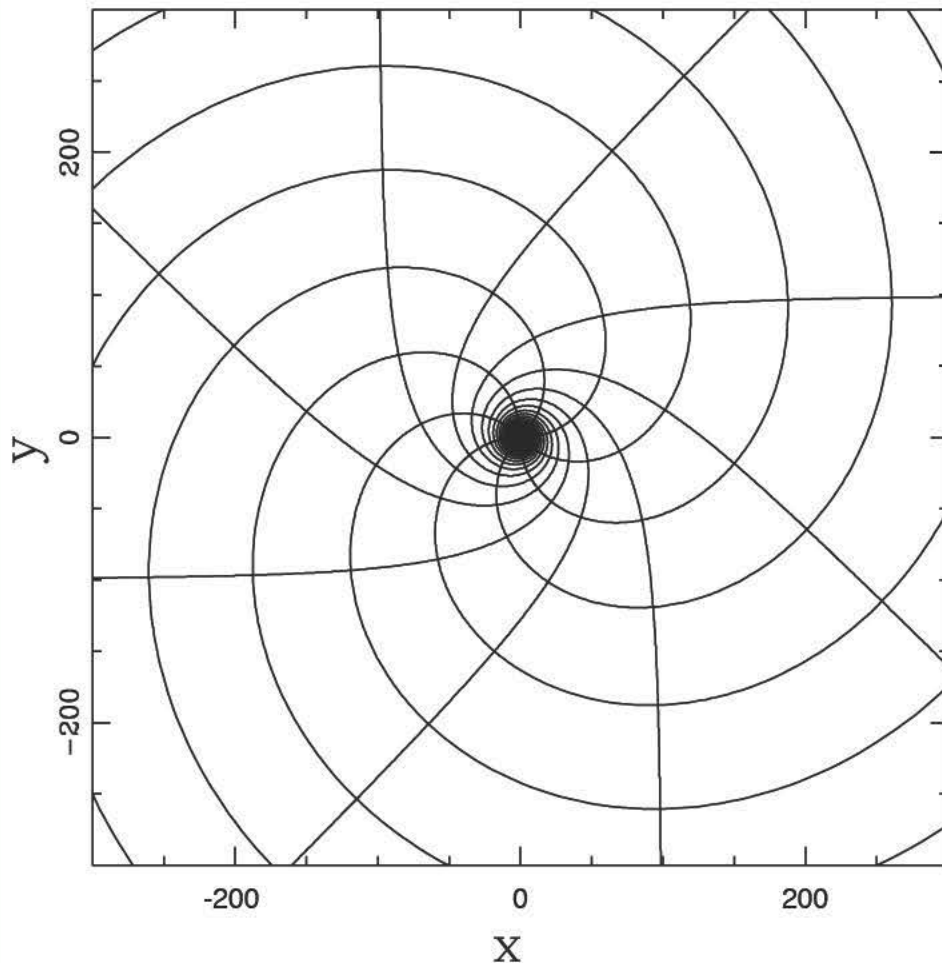
$$p = -\left(a^2 + \xi^2 + 2a\xi\cos\theta\right)^{-1/2}$$

$$q = \frac{\xi\sin\theta}{a + \xi\cos\theta}$$

$a = \text{radius of orbit}$

$$Y(\xi) = \frac{2\xi^2(1+q^2)}{a[\xi^2(1+q^2)\xi^2 - a^2q^2]^{1/2} + [(1+q^2) - a^2q^2]}$$

Parker Spiral in Equatorial Plane



$$p = A \log \left(1 - \frac{1}{\xi} \right) + \phi$$

$$q = \xi - 1 - \log \xi - A\phi$$

$$A \equiv V_{wind} / (\omega R_*) \approx 100$$

(A sets shape of spiral)

Summary

- Construction of coordinate systems (p,q)
- Generalizes to many astrophysical problems
- Can find sonic points and dimensionless mass accretion rates analytically
- Dipole + Octupole system: flow density (10x)larger, hot spot has higher temperature
- Magnetic truncation radius changes
- General constraint on steady transonic flow:

$$n > \ell + 3/2 \Rightarrow \textit{nearly isothermal}$$

(Adams & Gregory, 2012, ApJ, 744, 55; Adams, 2011, ApJ, 730, 27)

Dragons

- *This use of coordinate systems in this context only works for potential fields (no currents)*
- *The formalism has been developed for two-dimensional systems; can work for three-dimensional systems in principal, but complicated in practice*
- *Treatment (thus far) limited to steady-state (time independent) magnetic fields: magnetostatics not MHD*

