

Non-planet mechanisms for clearing and sculpting discs

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Hope to convince you:

- (Xray) photoevaporation now well understood
- Significant factor in disc evolution
- (Probably) responsible for ultimate clear out of protoplanetary discs
- En route, produces structures with properties overlapping those produced by planets



Are observed structures (gaps, holes) pure dust phenomena?

Probably not....

- Sharp edge due to radiation pressure?

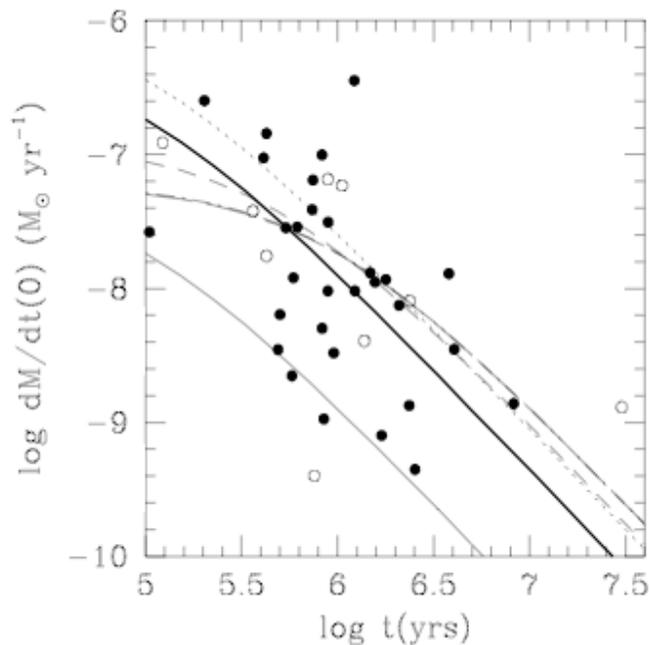
See Dominik & Dullemond 2011

- Sharp edge due to photophoresis (Krauss et al 2007)?

Hard to suppress small dust production

Do discs clear via viscous accretion?

- UV excess => do accrete at rates $\sim M_{\text{disc}}/\text{age}$ ✓



(Lynden-Bell & Pringle 1974,
Hartmann et al 1998)

Phenomenological description as due to action of (pseudo) viscosity e.g for $\nu \sim R$ ($\Rightarrow \Sigma \sim 1/R$)
get similarity solution:

$$M = M_{\text{in}} \left(1 + t / t_{\text{vis}} \right)^{-1.5}$$

- Far too slow decline at late times....

Need extra gas clearing mechanism

To clear discs within ~ 5 Myr, to make observed hole/gap structures...

- Planets ?
- Photoevaporation ?
- Clearing by MRI driven winds?

Suzuki et al 2010: inner hole forms very early (0.1 A.U. @ 10^5 years,
1 A.U. @ 10^6 years)

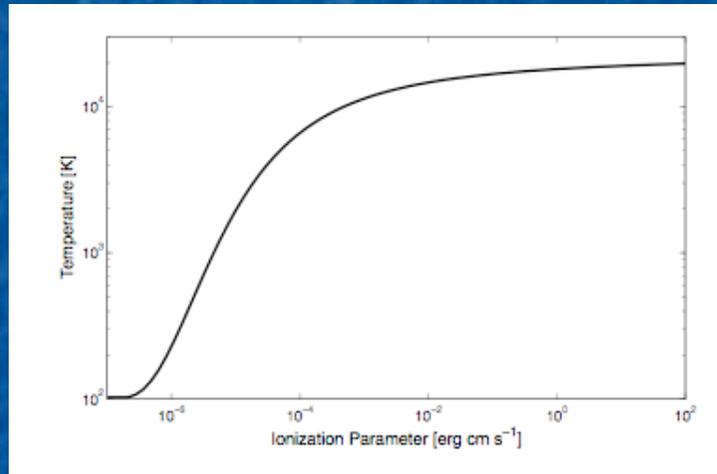
← Halts accretion onto star from very early times also....

Actually destination of 'wind' unclear (doesn't attain escape velocity....)

Key facts for understanding Xray photoevaporation

- Temperature of Xray heated gas set by

$$\xi = L_X / nr^2,$$



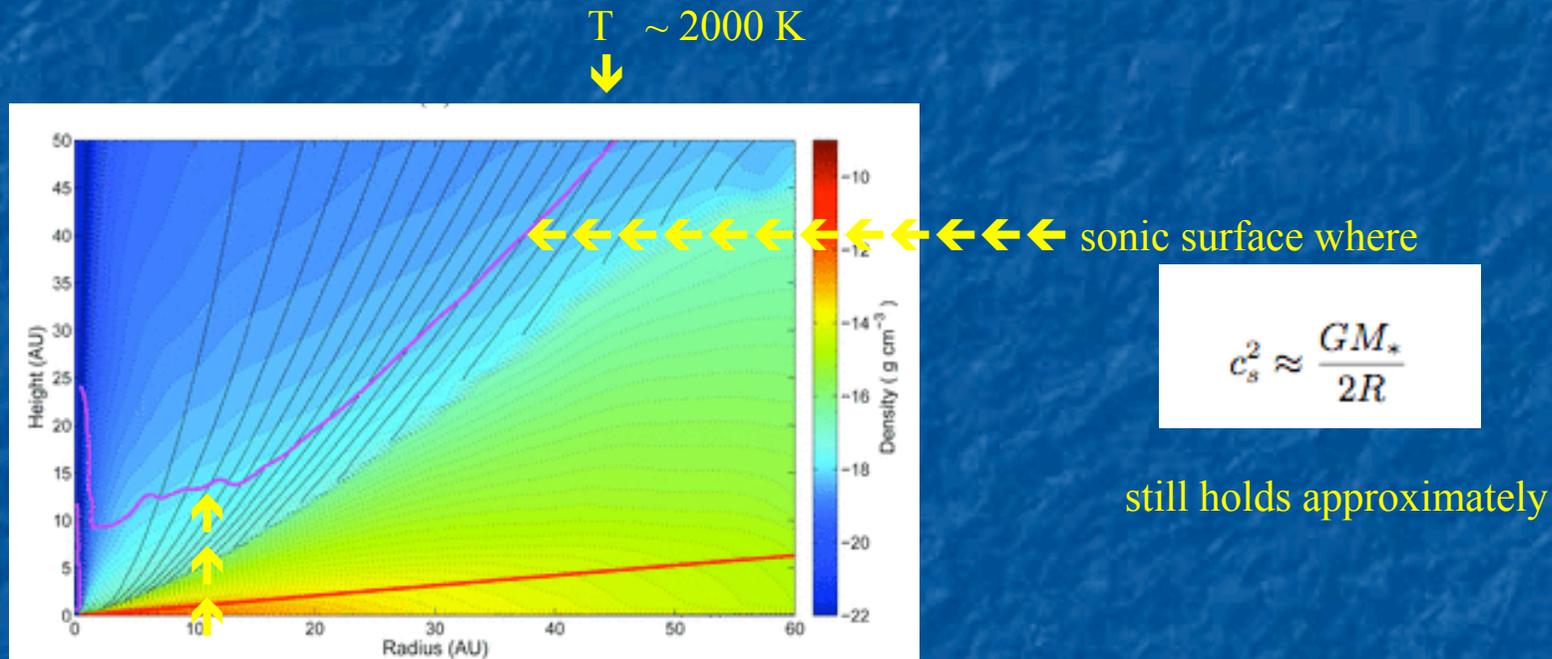
←←←← max T is 1-2 x 10⁴ K

- In a Parker [⊛]wind, sonic transition occurs at radius where

$$c_s^2 \approx \frac{GM_*}{2R}$$

⊛ = isothermal , spherical

Results of radiation hydrodynamical modeling of Xray photoevaporation (Owen et al 2010, 2011b)



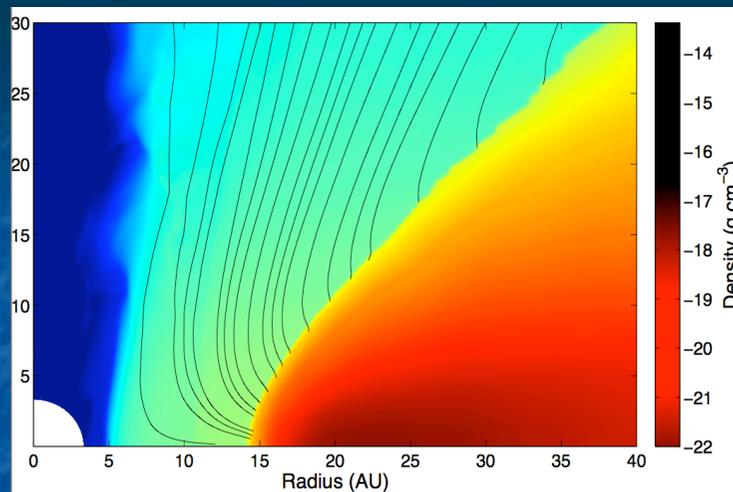
- $R \Rightarrow T$ and (beyond 10 A.U.) $\xi = L_X / nr^2$, fixes n and hence mass flux

Result:

$$\begin{aligned}\dot{M}_w(M_*, L_X) &= L_X \sqrt{\frac{8k_B\pi^2}{\mu m_h}} \int_0^\infty dT_{\text{esc}} \frac{T_{\text{esc}}^{-1/2}}{f^{-1}(T_{\text{esc}}/2)} \\ &\approx 8 \times 10^{-9} \left(\frac{L_X}{1 \times 10^{30} \text{ erg s}^{-1}} \right) M_\odot \text{ yr}^{-1}\end{aligned}$$

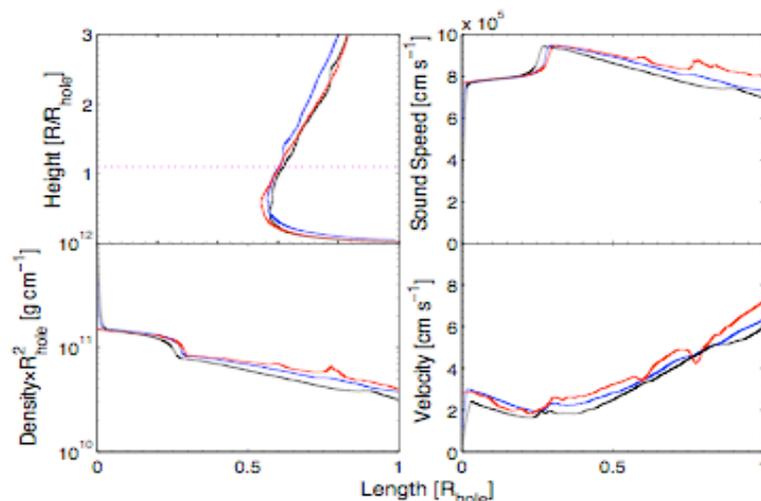
- Proportional to L_X
- Independent of M_*
- Doesn't depend on properties of underlying disc!

Discs with inner holes



Owen et al 2010

As vary R_{hole} , topology of innermost streamline and variation of c_s and u with scaled distance along streamline is invariant

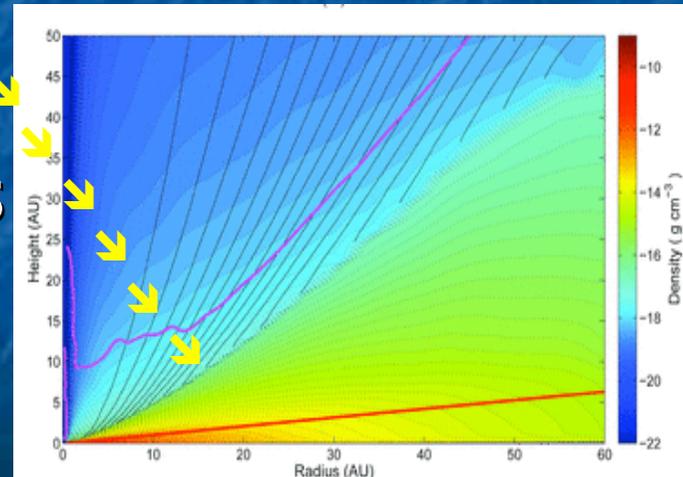
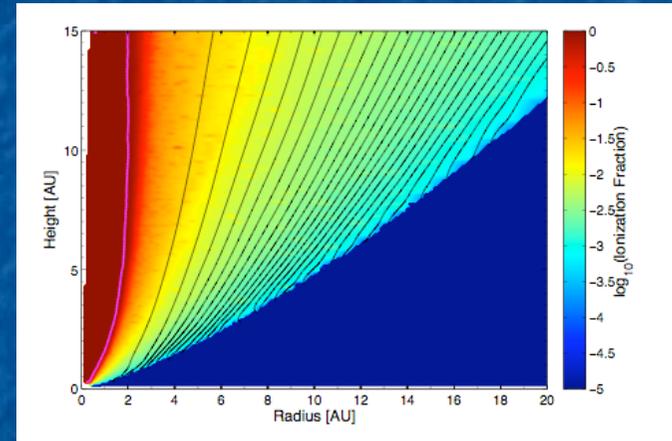


====> photoevaporation rate INDEPENDENT of inner hole size ($\sim 10^{-8}$ solar mass/yr for $L_X - 10^{30}$ erg/s)

Owen et al 2011b)

What about other radiation sources?

- EUV? Can't penetrate Xray wind
- FUV? Within 100 A.U. only heats below Xray sonic surface - doesn't change mass loss rates



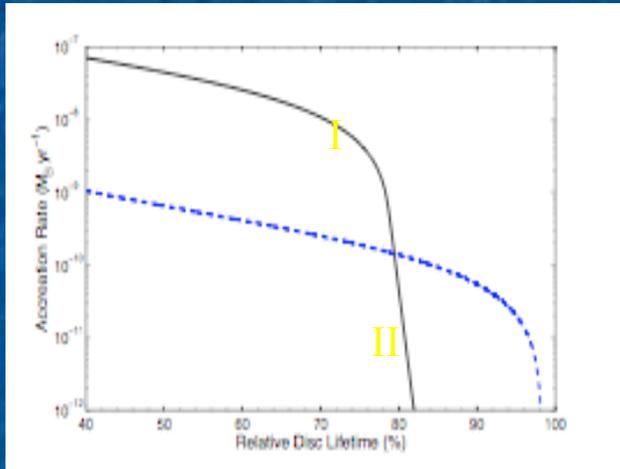
Owen et al
2011 b)

(But may affect structure of subsonic region:
See Gorti & Hollenbach 2004,2008,2009)

ALSO MAY BE IMPORTANT MASS LOSS MECHANISM AT > 100 A.U.

Combining photoevaporation with viscous evolution:

↖ Constant L_X



Initial →

75% total lifetime →

76% total lifetime →

77% total lifetime →

78 % total lifetime →

79 % total lifetime →

Etc.

Four stage evolution:

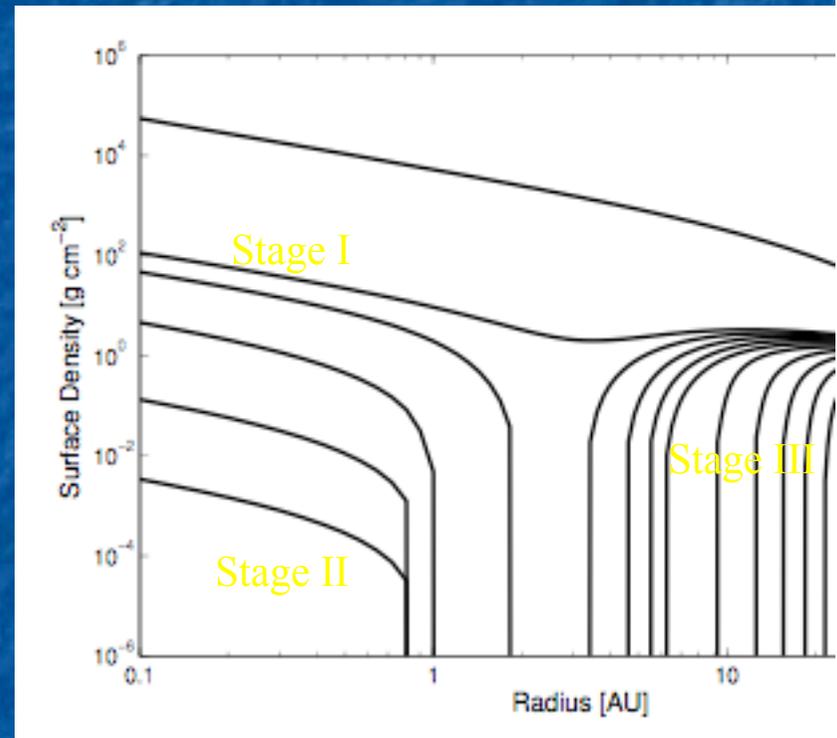
- I Viscous dominated
- II Draining inner hole
- III Outer disc clearing
- IV Thermal sweeping

← Accreting, dust poor (migration), <10 AU

← Empty inner hole, > 10 AU

← NEW

Owen et al 2011a)



↑
76 %
total
lifetime

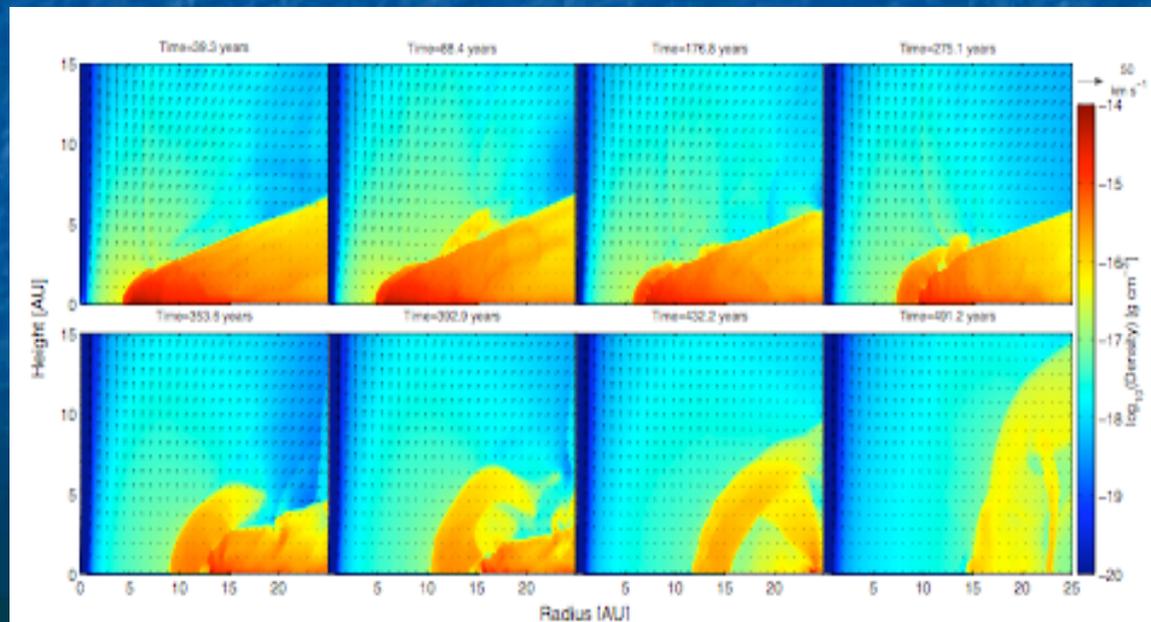
↑
80%

Stage IV: thermal sweeping

- Once Xrays penetrate a radial distance $\sim H$ into disc, heated gas evaporates vertically in 'plume flow'
- Residual disc clears on \sim dynamical time of inner rim (~ 10 s of A.U.)

Sets in when column density at inner rim is $\sim 0.5 \text{ g/cm}^2$

Remove few \rightarrow 10 Jupiter masses of gas



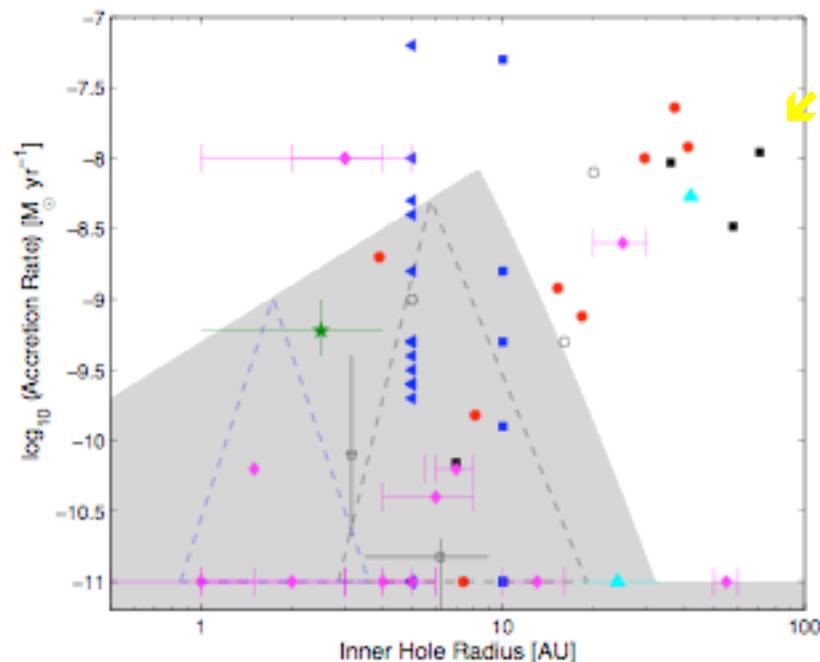
Thermal sweeping limits lifetime of non-accreting hole stage (stage III)

- Fraction of lifetime spent with hole (stage II + III) $\sim 10\%$
- Fraction of lifetime spent with 'transparent accreting' hole (stage II) $\sim 5\%$
- Fraction of lifetime spent with non-accreting hole (stage III) $\sim 5\%$

Which inner hole sources could be due to photoevaporation?

Around half (those in shaded region)

Owen et al 2011b



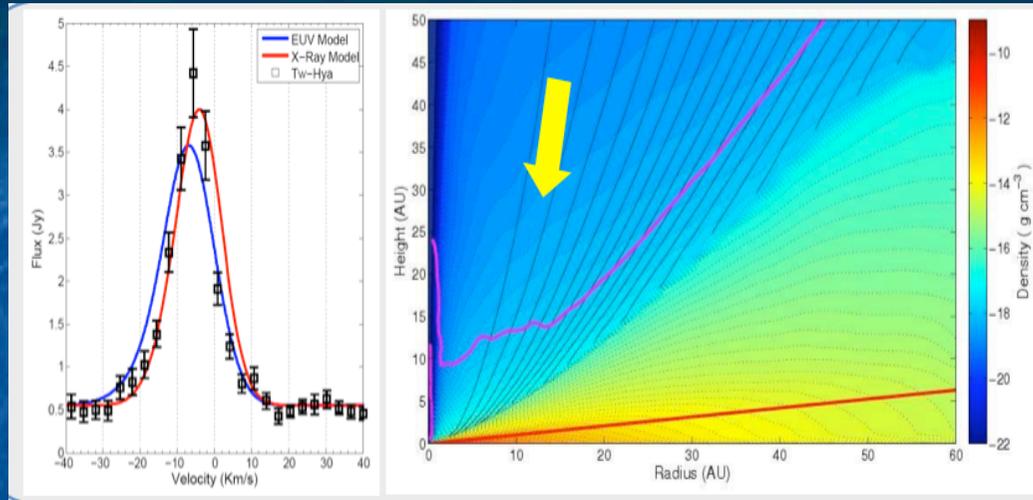
These can't

These are upper limits

- Systems evolve
 ↘ as inner holes drain
- Initial \dot{M} depends on L_X
- Initial radius depends on M_*

Cyan = Brown et al 09, blue = Cieza et al 10,
Black open - Ercolano et al 09, Black filled =
Espaillat et al 08,09, red= Kim et al 09, magenta -
Merin et al 10, green = Najita et al 10

Evidence for Xray photoevaporation



Xray: Ercolano & Owen 2010
EUV: Alexander 2008

- Both Xray and EUV photoevaporation explain line profiles of NeII 12.8 μm
Cf observed profiles for TW Hydra, Pascucci & Sterzik 2009
Pascucci et al 2011
- Only Xray photoevaporation explains low velocity ($\sim 5 \text{ km/s}$) component of OI 6300 in T Tauri stars (cf EUV models: Font et al 2004)

.....but note lack of blueshifted OI 6300 in TW HydraPascucci et al 2011

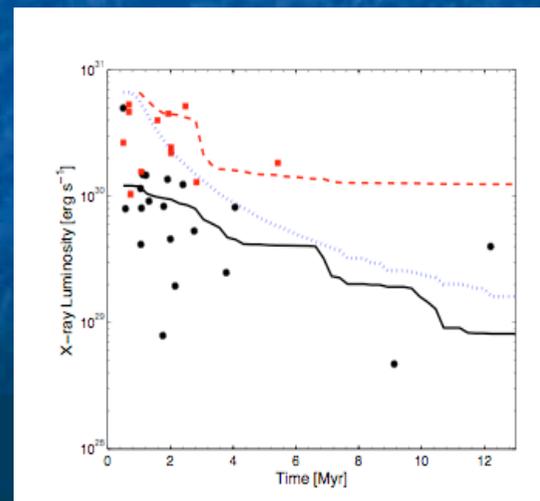
Evidence for Xray photoevaporation?

- High L_X stars lose discs earlier - implies WTTs should have higher L_X on average.

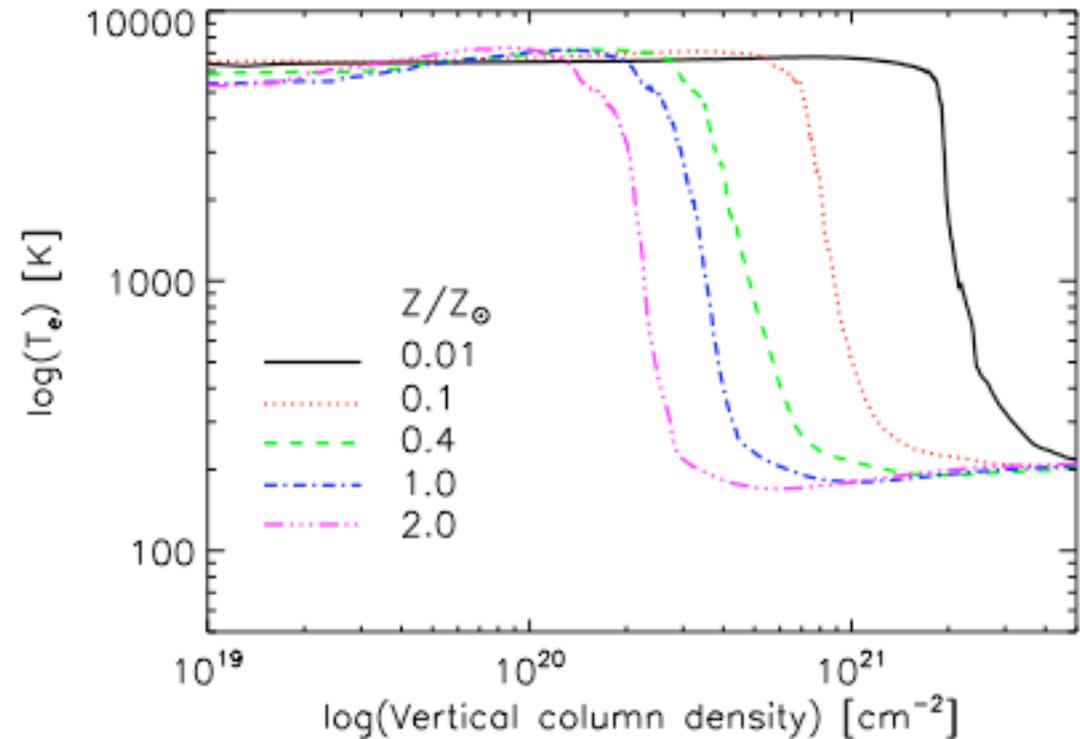
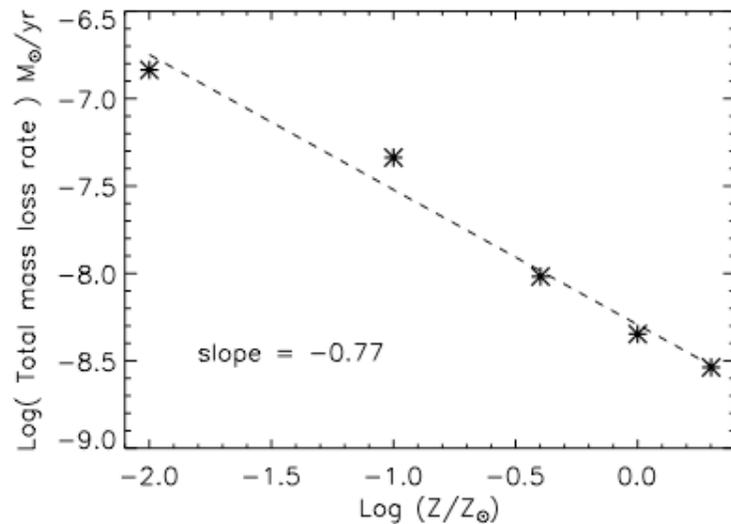
Preibisch et al 2005,
Gregory et al 2007

Well known observational correlation. Usually argued that Xrays suppressed/absorbed by accretion: perhaps instead accretion suppressed by Xrays.....

Population synthesis: Owen et al 2011a



Dependence on metallicity:



Photoevaporation more efficient at low Z: lower dust extinction => Xrays heat to higher column

EXPECT SHORTER DISC LIFETIMES AT LOW Z

Disc lifetime ****increases***
 strongly with decreasing Z
 if it's instead set by time
 required for planet formation

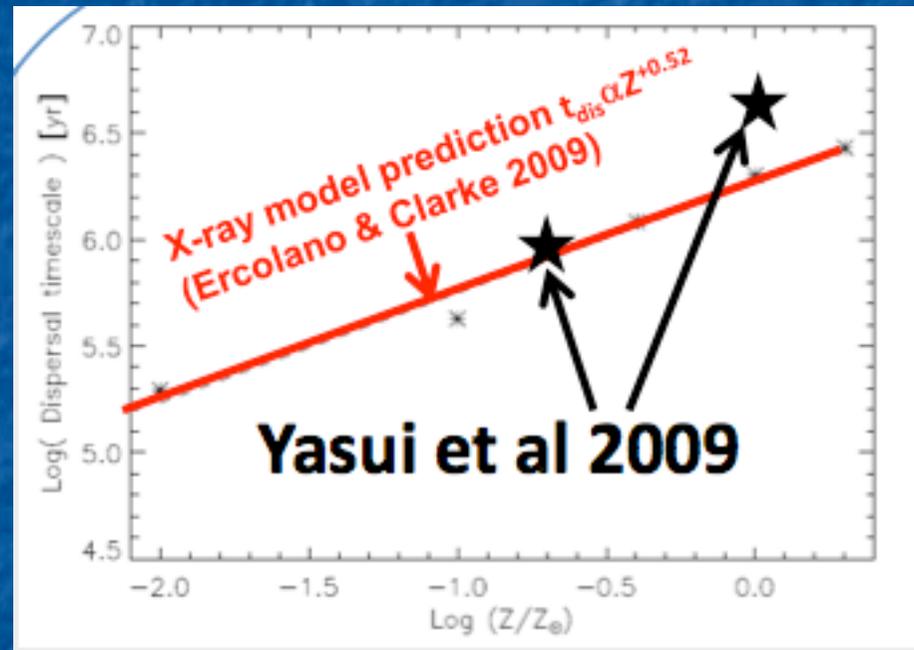
$$t_{form_{min}} \propto \left(Z^{-5} M_{g0}^{-7} r_{d0}^{(9-5p)} \right)^{(2-p)/(5-3p)} \propto Z^{-5/2}$$

(see Ercolano & Clarke 2009)

$$t_{form_{min}} \propto \left(Z^{(-13+2p)} M_{g0}^{-17} r_{d0}^{(4-2p)(6+p)} \right)^{1/(5-3p)} \propto Z^{-11/2}$$

- A possible observational discriminant?

- Recent claim of shorter disc lifetimes in lower Z environment



...further studies at low Z may hold the key to discriminating between photoevaporation and planet formation

CONCLUSIONS

Xrays can drive photoevaporative winds of 10^{-8} M_{sun}/yr at upper end of XLF: like EUV winds, these produce a RAPID clearing phase but the Xray wind cuts in at much higher accretion rate.

Produce small holes at range of accretion rates but no accreting holes beyond ~ 20 A.U.; expect accreting and non-accreting holes to have similar frequency.

- Xray photoevaporation \implies line diagnostics ([Ne II] 12.8 μm and [OI] 6300 lines)
- Xray photoevaporation: shorter disc lifetimes at low Z (opposite to clearing by planet formation)

